

requires! 'le distinguish de rio inhretises influences r{ui paraissent intei'venir et dont les actions, coi "1' r{ii('es des orientations, se combinent récipro- fJtic men t: influen cc direct stu' the corner not }'ar a kind of 'rubbing or d'en traî même ri t et in11 ucnc'- sur l'ai iriantiitioii }x op'i'e bles ;ii quilles due au déplacern ent in pidc diins le c)iam J , sans li:irlcr des ccii rati ts engendrés pro bablcment dans l'i ctt vc r}u i j' ue en incrnc te mps tin rûle d'(Ecran and can partly mask these ell'ets.

Mais, cont iiic- on le conçoit faci le mien t, tlans bles conditions telles c}ue celles où nous opérions (soit 'la ns un wagon .d'u ri train rapide ou sur une machine isolée) et qui ue }icrine tien t de fii ire toutes les manœuvres désiratiles, par suite dii mia{; n/tisiie sous-}ieriiianent dcx elo}'pè }iar les vilii "utions et the router in1, .iirif iue of the installation where we are at high speeds, it was not }iossible to detei'iiiinate the diveiv-es magne- data.

t'Jues and even less to dissociate ct d'É'tti il ier these different actions. Ces sans doute nc se iii1'len t }'as i 'csenter uti int(r^ l' itir}ti* iiiii mt- diat, mais elles oc sont toute l'ois poin t sans port ce ggenerale, et nous mous propose rl'cn continue lii rechei'clic.

O PTIQUE. - *Les syslénies optique.s en rnoireritrnf et la ircinslalioii Ille la Ten e.*

Note by Mr. G. Saosac, presented by Mr. Lippmann.

1. Efrt fr rlémeittaii-e movement. - J'ai expliqué einématijuement : l'enti.i i ncuent [Partiel des ondes luini iieuses pai' l'eau en niouveine ri t *Comptes ren lus*, t. I 2ñ, p. 8 i 8; *Soci'éld fi'ançaise de NÂJ sis-*^ 'o99) i = Yeltnann's principle and astronomical aberration studied with a system optique t{uelcont}uc *Comptes rencliis*, t. 141 , I go5, p. i cao). Mes raisonnements supposent que l'etlicr du vide n'est pas entraîné tout dans la traduction de la matière (hypothèse de Fresnel), ou, dii moins, que la vitesse r du s3-stème -rllique }i.ir rapport fi l'ether du vide est uniforme aii.x points of the system. But whatever the distribution of the vector v in the scope of the system, it is periiis to keep in the following form the principle of the cllementary motion effect, which I established in 1°99 (toc. cit.) and which will serve as a breeze for a more cineniatic theory. genéi'ale.

Sur char}ue elê-nien t de loiig'ucur *cll* lie à u ri système optique, la trans- lation du systeme fait varier la durée de propag ation des ondu lations lumi- neuses de ^u *dl* (effet dc \Mouvement eleiræn taii e); u désig'ne la composante,

according to 'll, of the velocity v of the element Cf of the system with respect to the vacuum ether; V, designates the velocity of light '*Inns the tyree, iT1*' in the element *ch* is coin}aris clans l'un des rn il ieur matériels du système atirjiic.

2. *Efel. toni htlonnaire maligne.* ΔT $\frac{C}{V_0^2}$ ainsi la variation AT que la dur ec de propagation sur le périmr-tre du circtit subit sous.s l'influence du relative movement cle this invariable circuit and ether 'lii vide. This is the sum of clem on taii'es Vé c tendue fi tons les elena ents 'll alu circuit. Or, la somme des aletit s Ile rr r/r-1-r:sente (lorcl I(e)vin) la r:rro/"iir'n C cle l'e ther le lon g du circeti on (i3jerl'nes) 1/n/erzsiffr ':*Jii lourl'i'llon* corres- }iondan t, through the circuit. Iii ti'ofoissons la valeur ur moyer ne / d u vecteur de Bjerl'nes, ou densité tlu tourbillon, p_- l'endiciilaireinent à la siii'face S dli Cil'CUi StlppOST }11 éfl . J-C fPt tOtl'l')J11O11fililrC Optit{tle 8 }JOUI! Villett11'

$$(1) \quad \Delta T = \frac{C}{V_0^2} = \frac{b S}{V_0^2}.$$

If the density of the tourlaillon is toijotir.s ri i:lle, in other words, if the relative moirre- iiient tle l'etlier is *irrolalu'itnrl*, the valcNr of AT is zero and we }ieu t apply Vel tinan ri's theorem ('c. "i").

If, on the other hand, the relative slack is *rotated nell* the delay dT {" ocliit a variation 'J- 1° "ise (L, long'iierir d'onrle) :

$$(2) \quad x = \frac{b S}{\lambda V_0}.$$

then interferer bleu x systems of light ondiilations '{ii i have traversed in opposite directions the circuit o} tir}ue Ile large surface S (see inès

Notes, *Comptes i eii':lus*, t. 160, I t) 1 O 1° $\wedge 0^*$ $\wedge 0^* 17$)

L'effet tourlaillonnaire alterera de zm 1'i difference Ile phase des deux ondulations inverses, car il rcsul te rl'effets de mouvement elu }" einier ordre c}ui cliaig'ent Ile sens avec la }= 1"3'ation Ile la lumière.

3. *Limite siipérview-e Ile l'r.nnt'cn "ienienl rl,e l'elliei- rlcins In trcmslnli'oii 'ie In Tew'e.* - Si l'étlier est su'}iose entraîné aii oisina ge du sol, la vite.sse relative r de la Terre et clv 1'et lier aug'iiiente de fr r{uanrl l'altitude croit 'le As et ne devient ég'ale 'i lit vitesse r, de ti.iii.slation Ile 1:i Terl'e r}u'à l'altitude où cesse ï'entrainenien t.

Around noon (or midnight), the speed v is Jaarall!ele at the horizon, the vector ô is horizontal, close to the inericlian, and has the value ps if we

neglect the

curvature of ether flux lines with respect to - Under these conditions, the value (*) of m applies to the surface S of a vertical circuit oriented east-west.

From noon to midnight, the direction of propagation of each ripple is reversed in space, the 2s variation of the phase difference is inverted and the interference bands must move by $d m$ rows.

In the course of observations that I will describe elsewhere, I found that the position of the central fringe of my 'i inverse beam interferometer (*loc. cit.*, {i. i G26) did not depend on the time of day. my pointing accuracy made it possible to determine an upper limit of m corresponding to, "" of — wavelength for a circuit of 30° contour, inclined to the horizon; vertical projection zo''' . According to the formula (*), \dot{u} or $\frac{\Delta r}{\Delta z}$ admits the radian per second. In other words, for a vertical ascent of I'' , the relative velocity v increases by no more than a fraction $\frac{1}{3} i o''$ of the Earth's velocity v .

Taking up the theory of star alignment *Comptes rendus*, '90°. *loc. rit.*) in the liy}iotliuse of an entrainment of the etlier r' us the ground, one sees easily c{u'elle subsiste, à condition de définir l'aberration par la vitesse

relati -e r of the globe and the étlier att *place of observation*. As the value of the Earth's velocity v , makes the oliserved value of the alierration retron ver, to the a}rroxiiriation of , is that the entrainment $v e l o c i t y$ ($r, -v$) near the ground adrriet Jr, as the su}ierieure limit. The result of these observations completes the r-ecedent. Moreover, it shows c}u il fait réduire lieaucou}i la upper limit lle the drive speed ($r - r$) if you do not want to This speed is still noticeable at great altitudes.

6. *Effet toui-hillounnire opticJue a, agiilaire.* - Let two Jinettes be directed towards each other at a great distance D . On the area ($D . /$) of the section (diameter f) of the long beam l unieneu.x which separates the glasses,

l'ef't et tourlillonnaire $1' d$ nit le i et'ird ET ou $D . /$ en entre les deux vilx a-tions élénientaires p- iagées suivant les l'ordS $1^\circ 1^\circ$ Osce du faisceau. To tue the focal synclironisine is ret,abli, the image of the foyet' of a telescope in the focus lle l'autre 'loit é tre deviée tl'ti n any'lc t tel '}ie l'avance géométrique c1 compensates for the geoinetric}ue reta rd V, ET. It is easy to deduce that if the sharp lii $L_$ is exactly pointed at the lunet te L , , the latter is depointé-. $1'-rP'$

'I " of angle $2t$ or $\frac{AD}{V}$

According to the upper literature of \hat{u} , c}ie inès observations ont ('table, the effect

tonbillonnaire angulaire * e admet la limite supérieure $\frac{3}{c} \text{ i o}'$ -

To determine this limit directly, it would have been necessary to set the precision of the reciprocal tips of the two glasses at less than o'' , through a mesoplieric layer $i oo'$ long, or even at less than o'' , $i' a$

15° Island distance.

TH ER MODYùTAMIQU E. - *A pylicnl.i'on dif I*'-*-cij'e Ile heu - cmæ ju/iènornenes yti accomyamet lci charge des cozdrtslnleurs.* Note by il. A. EDLti, presented by II. E. Bouthy.

I will first summarize, in a single section, the demonstrations relating to . t these

1 lic-nomena given by Pellat (') and by Sacerdote (') using Massieu's remark.

Soi $t u n co nrfenscifeii$ / *fer mē* don t les a i in atit res in $i n$ ces son t coll'ies su t le d ielet- trique. L'a rma $t u re ex te rn e B c-onam u ri iq ue où l'e ricei n te, ta n dii q u e l'i ri tern e est } iortce a u pot en t tel $V > o$.$

The state of this condensate $t1/$ depends on V , its tent perat u re T'' and 'the p ressi on s u ii i-forms P and p qu i ni g n e, one at the extu rien r, the other at the i ri tc ri eu r.

Let's assume 1° co nstant.

Sous $1'i ri fl uen ce des vari al ions d V, d'1', d'', la clia rie 31 de \lambda a u, in e ri te de rf\lambda f, le dié lectri c que reçoit une qua ri ti té de cli ale u r/Q, le s ol u nde in tc ri cu r (ca vité) aux ir e ri te de dv e le vol u ir e ex rérie ur v' de di''$.

The energy increase o f the condensate in this process is

$$(1) \quad dU = V dM + J dQ + p dv - P dv'.$$

If this transformation takes place in a way that is similar to il' le, rfU, $d(SIY)$, rt(/ai) and c/ ($1^\circ i''$) are d i Her en tielles ex actes, and it is Ile iué in e de

$$(2) \quad dX = dU - d(MV) - d(pv) + d(Pv') = - M dV + J dQ - v dp.$$

Poso ns

$$(3) \quad dQ = a dV + b dp + c dT$$

(1) PELLAT, *Journal de Physique*, 3^e série, t. VII, p. 18.

(') I*. P'r.ÉtuoTe, ' lit.e cle rloclora I, y. fi. Pa i'sis. i Sgg.